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Cover Photo:

The cover photo shows two crystals to be used for the CRESST (Cryogenic Rare Event Search with Superconducting Thermometers) experiment located at the Laboratori Nazionali del Gran Sasso (LNGS) in Italy. The CRESST experiment is designed for direct detection of weakly interacting massive particles, accounting for the dark matter observed in the universe. The Max-Planck-Institut für Physik is substantially contributing to all essential aspects of the experiment.

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Chapter 1

The ATLAS Experiment at the Large Hadron Collider LHC

1.1 The Muon Spectrometer

The muon spectrometer [1, 2] of the ATLAS experiment is equipped with three layers of muon detectors in a toroidal magnetic field of 3 – 6 Tm bending power generated by a superconducting air-core magnet system. The spectrometer is designed to provide muon momentum resolution of better than 10% for transverse momenta up to 1 TeV/c over a pseudo-rapidity range of $|\eta| \leq 2.7$. This requires a very accurate track sagitta measurement with three layers of muon detectors which have to be aligned relative to each other with an accuracy of up to $30 \mu\text{m}$ in the bending direction in the magnetic field. Drift chambers with very high spatial resolution of $40 \mu\text{m}$, the Monitored Drift Tube (MDT) chambers, have been developed to cover the active area of the spectrometer of 5500 m^2 with only 5% gaps mainly in the region of the detector feet.

The cylindrical central part of the spectrometer (barrel) contains eight race-track shaped magnet coils of 25 m length and 5 m radial width. The layout of the muon chambers follows the eightfold symmetry of the magnet around the proton beam axis in eight small and eight large azimuthal sectors (see Fig. 1.1). The barrel part of the spectrometer is complemented by two endcaps each consisting of eight superconducting coils housed in a common cryostat fitting into the inner bore of the barrel toroid magnet and three wheel-shaped layers of muon detectors.

The MDT chambers built in Munich consist of two triple layers of 30 mm diameter aluminium drift tubes of 3.8 m length equipped with a central gold-plated tungsten-rhenium sense wire which are separated by an aluminum space frame. The drift tubes are operated with Ar:CO₂ (93:7) gas mixture at a pressure of 3 bar and a gas amplification of 20,000 (corresponding to a operating voltage between tube wall and wire

of 3080 V) and provide a position resolution of about $80 \mu\text{m}$. The sense wires are positioned within a chamber with an accuracy of better than $20 \mu\text{m}$ in order to achieve the required spatial resolution of the chambers [3].

In the years 2001 to 2006, 88 MDT chambers were constructed at the MPI [4] containing about 36000 drift tubes of 3.8 m length. They cover about 15% of the active area of the spectrometer. The installation of these chambers in the outermost layer of the barrel part of the spectrometer took place in 2005 and 2006. The commissioning of the chambers in the ATLAS detector was completed in summer 2008, just before the start of LHC data taking [5, 6, 7]. The one year long interruption of the LHC operation caused by severe beam magnet failures in September 2008 was used to calibrate and align the muon spectrometer and to commission the muon identification with muons from cosmic rays.

1.1.1 Muon Detector Software

The MPI group is contributing to all aspects of the software development for the analysis of the ATLAS muon detector data, including the calibration of the space drift-time relationship of the MDT chambers [8, 9, 10, 11, 12, 13, 14, 15], the alignment of the muon spectrometer with muon tracks [17, 18, 19, 20, 21], the evaluation and monitoring of the data quality of the muon detector [22, 23], optimization of the muon reconstruction for the high background rates expected at the LHC [19, 24], and the Monte Carlo simulation programs for the muon spectrometer [25, 26]. The first two projects which are of central importance for the operation and performance of the ATLAS muon spectrometer and in which the MPI group plays a leading role are discussed in more detail in the following.

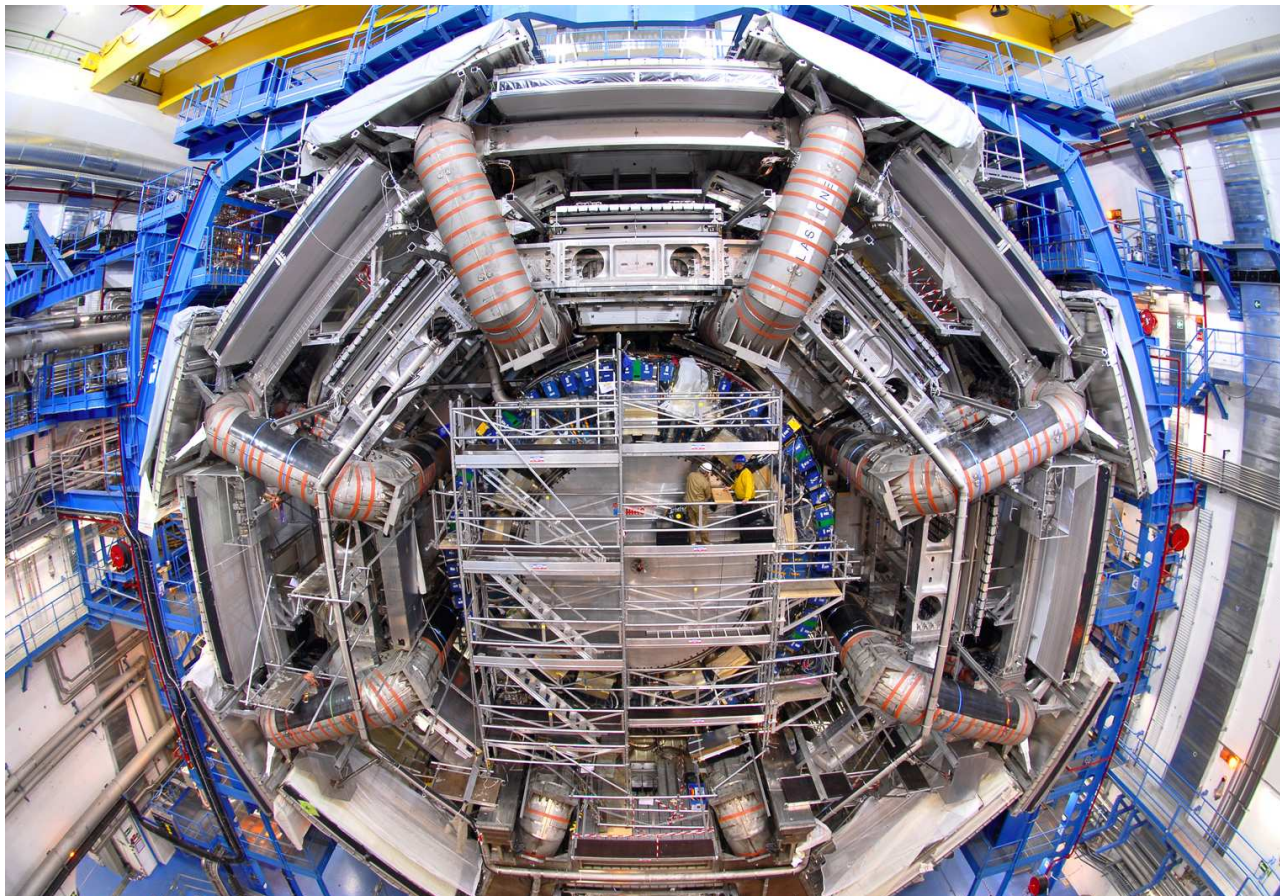


Figure 1.1: The barrel part of the ATLAS muon spectrometer viewed from the LHC tunnel after completion of the muon chamber installation in November 2006. The MDT chambers built in Munich are mounted on rails on the outside of the eight superconducting magnet coils.

Drift Tube Calibration

The pressurized drift-tube technology has been chosen for the ATLAS muon tracking chambers because of the high spatial resolution of better than $80\ \mu\text{m}$ which can be achieved for individual drift tubes by measuring the drift time of the ionization electrons to the sense wires. In order to reach the required spatial resolution of the chambers of $40\ \mu\text{m}$ not only the sense wires have to be positioned in a chamber with an accuracy of $20\ \mu\text{m}$ but also the space drift-time relationship of the approximately 400,000 drift tubes in ATLAS has to be known with the same precision. The drift properties of the electrons depend on the temperature and pressure of the drift gas as well as on the magnetic field and the rate of background hits from neutrons and photons which can be very high at the Large Hadron Collider (LHC). For the regular calibration of the space drift-time relationship, correlations between the positions measurements of the drift tubes hit by a traversing muon can be used.

We have developed fast and efficient algorithms for the determination of the space drift-time relationship [8, 9] and of the spatial resolution [10] of the MDT chambers which are now part of the ATLAS data reconstruction software. The algorithms have been extensively tested with simulated data, cosmic ray commissioning data [8, 6, 7], in a muon beam at CERN [11, 12] and under LHC operating conditions, in magnetic fields and at high background rates, using test beam data taken by our group at the Gamma Irradiation Facility at CERN [13]. A model for magnetic field corrections to the space drift-time relationship has been derived from the measurements [14].

The large sample of cosmic ray muons recorded by the ATLAS detector in 2008 and 2009 allowed detailed studies of the robustness of the calibration algorithms. The r - t calibration algorithm showed robust operation for all the chambers apart from the chambers with muons at 30° incidence angles where all drift radii are equal. The r - t calibration method originally applied

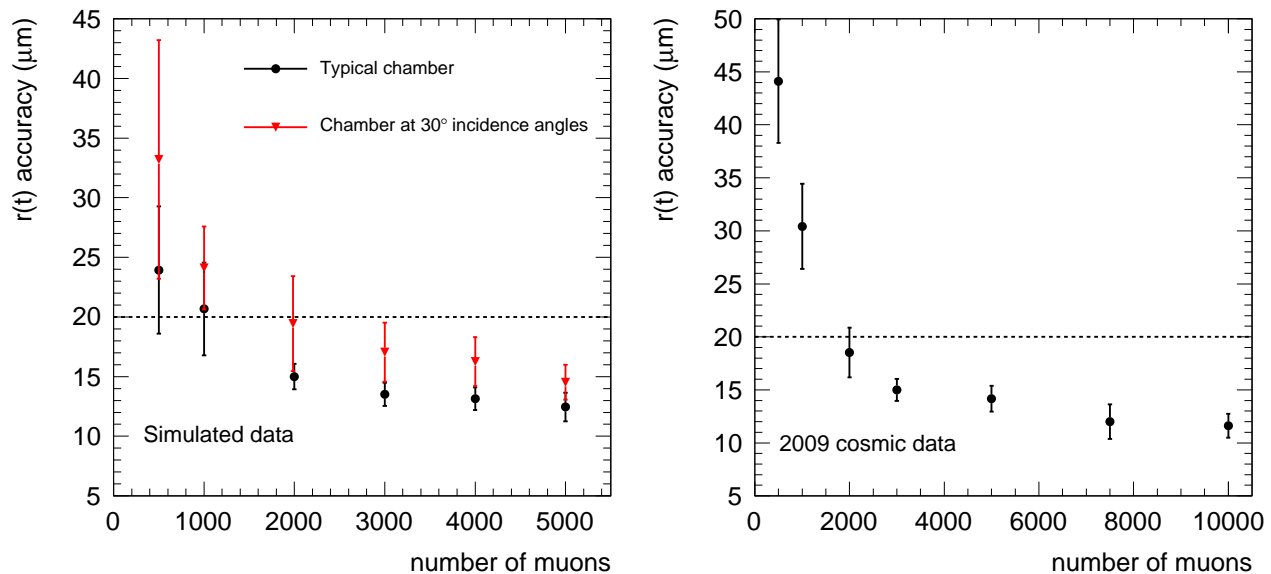


Figure 1.2: Accuracy of the calibration of the space drift-time relationship of the drift tubes of the ATLAS MDT chambers as a function of the number of cosmic muon tracks used per chamber. Left: Simulated data. Right: Cosmic ray data confirming the predicted r - t accuracies.

to straight track segments in triplelayers of the muon chambers had to be extended to curved track segments in the muon chambers to achieve the same level of robustness for all muon chambers [27]. Figure 1.2 shows the accuracy of the space drift-time relationship provided by the improved calibration algorithm as a function of the number of collected muon track segments in a chamber for simulation and cosmic ray data. The required accuracy of 20 μm is reliably achieved with 3000 muons per chamber for all chambers of the muon spectrometer. The predicted accuracy is confirmed by the studies with muons from cosmic rays.

The dependence of the drift time on the magnetic field strength measured with cosmic rays in the ATLAS detector agrees well with the model based on the test beam measurements (see Fig. 1.3).

The MPI group is operating one of the three computing centres dedicated to the calibration and the alignment of the ATLAS muon spectrometer [15] using a special data stream of muon tracks reconstructed by the ATLAS second level trigger algorithms. The commitment of the MPI group to calibrate roughly one third of the MDT chambers and to align the muon detectors with muon trajectories includes the development of fully automated calibration and alignment procedures and monitoring tools [16].

Muon Chamber Alignment with Tracks

To achieve the required momentum resolution of the ATLAS muon spectrometer up to the highest muon energies, the relative positions of the chambers have to be continuously monitored and misalignment corrections have to be applied to the measured track sagitta with an accuracy of 30 μm . The MPI group contributed significantly to the development and test of the high-precision optical alignment monitoring system for the muon spectrometer [1]. The group also developed algorithms for the alignment of the ATLAS muon spectrometer with muon tracks [17, 18, 19, 20, 21].

In the barrel part of the spectrometer, only the large chamber sectors mounted in between the magnet coils can be fully aligned with optical sensor measurements. The small chamber sectors mounted on the coils (see Fig. 1.1) have to be aligned with respect to the large sectors with muon tracks passing through the overlap regions between the the small and large sectors. The MPI group is providing these alignment corrections using the muon calibration data at the Munich calibration and alignment computing centre [15, 17].

Straight muon tracks measured while the toroid magnets are turned off are needed for the precise determination the initial chamber positions after installation as a starting point for the monitoring of further chamber movements by the optical sensors mounted on the chambers. An efficient algorithm has been de-

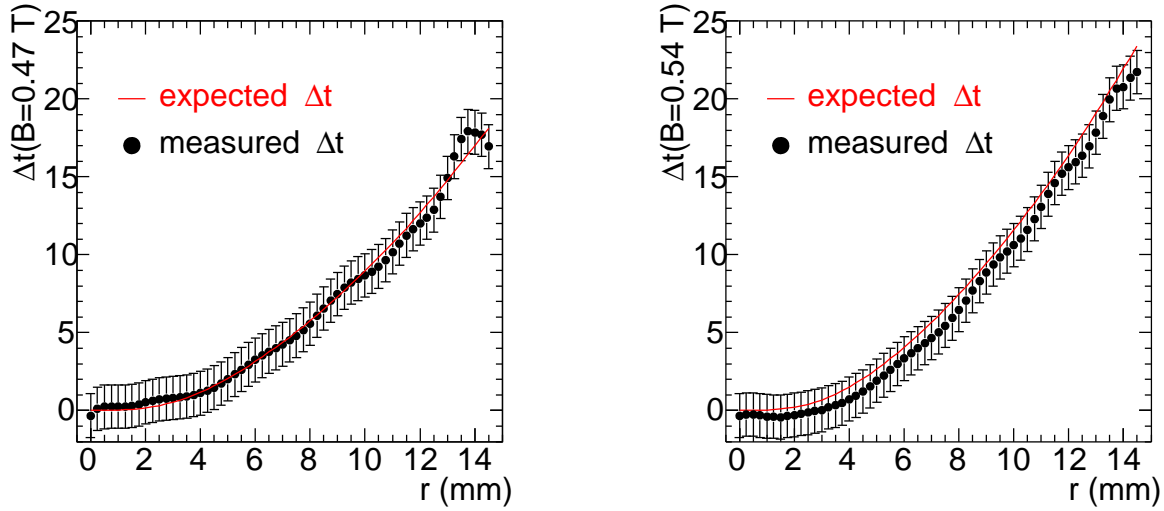


Figure 1.3: Magnetic field corrections to the drift time t in different regions of magnetic-field strength (0.47 T on the left, 0.54 T on the right) in MDT chambers in the middle layer of the large sectors of the barrel muon spectrometer as a function of the drift distance r . The measured difference in the drift times with and without magnetic field B agrees well with the model expectation in red derived from test beam measurements. The error bars of data points correspond to the estimated r - t accuracy of $20 \mu\text{m}$. The drift time correction increases with the distance to the wire because of the deflection of the drifting electrons in the magnetic field oriented in the direction of the tube axis.

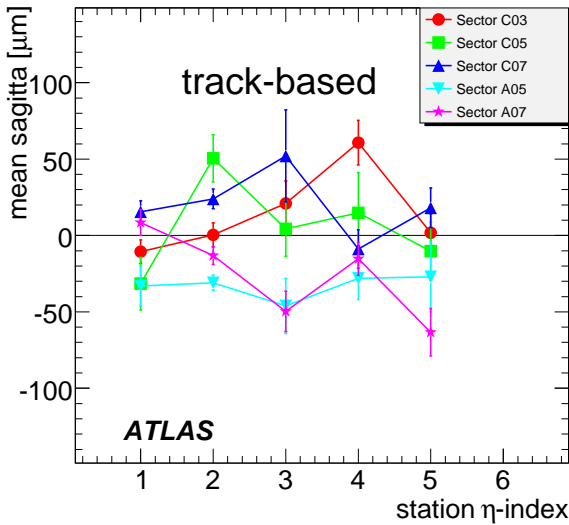


Figure 1.4: Mean values of the apparent sagittas of straight muon tracks after the initial alignment with straight muon track in ATLAS cosmic ray data. The station η -index labels the chamber position from the inside to the outside of the barrel of the muon spectrometer. The results are shown for large barrel chambers; similar results are obtained for the small barrel chambers.

developed and successfully applied to cosmic ray commissioning data [18]. Figure 1.4 shows the mean value of the apparent sagitta of straight cosmic muon tracks after the initial alignment with straight cosmic muon tracks. The accuracy of the initial alignment is of the order of $50 \mu\text{m}$ or better and close to the desired ul-

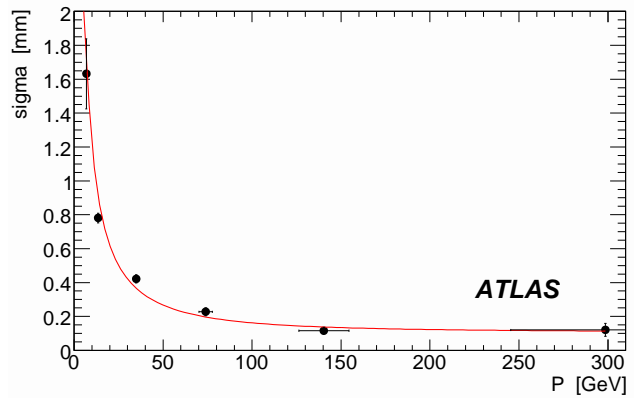


Figure 1.5: Resolution of the sagitta measurement as a function of the muon momentum as measured in ATLAS cosmic ray data. The red line illustrated the fitted resolution which is the quadratic sum of a term proportional to the inverse momentum taking into account multiple scattering and a constant term (of $100 \mu\text{m}$) reflecting the limitation of the resolution by the spatial resolution and the alignment of the muon chambers.

mate accuracy of $30 \mu\text{m}$. Improvements of the alignment accuracy are expected when chamber deformations will be taken into account in the track reconstruction. The measurement of the sagitta resolution of straight tracks in the muon spectrometer as a function of the muon momentum measured by the inner detector is presented in Figure 1.5. It improves with increasing momentum as multiple scattering decreases

with increasing momentum and reaches a plateau value of about $100 \mu\text{m}$ at high momenta reflecting the limited spatial resolution and the residual misalignment of the muon chambers. The achieved resolution is about a factor 2 larger than the target value of $50 \mu\text{m}$ because the chambers are only aligned with $50 \mu\text{m}$ instead of $30 \mu\text{m}$ precision.

The method is being extended to the alignment of the muon chambers with curved muon tracks in the magnetic field during normal operation of the experiment in order to verify the optical alignment corrections. This requires an independent measurement of the muon momentum which is insensitive to misalignment of the MDT chambers along the muon track. The MDT chambers with their two triple or quadruple layers of precisely positioned drift tubes measure not only track coordinates with high accuracy but also the local track direction. This feature can be utilized for independent momentum determination from measurement of the track deflection angles between the inner and the outer chamber layer and between the two multilayers of the chambers in the middle layer located inside the magnetic field. Studies with simulated data showed that the required alignment accuracy can be achieved within two days of data taking at the nominal LHC luminosity of $10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ [19, 21].

1.1.2 Measurement of the Performance of the Muon Identification in ATLAS

The huge sample of cosmic ray muons made it possible to study the performance of the muon identification in great detail up to muon momenta of $300 \text{ GeV}/c$ [28, 29]. Figure 1.6 shows the track reconstruction efficiency in the ATLAS muon spectrometer for cosmic ray muons reconstructed in the inner detector. The efficiency is close to 100% in the instrumented regions of the muon spectrometer. The drop of the efficiency at $|\eta| = 0$ is caused by the acceptance gap of the muon spectrometer which is needed for the services of the inner detector and the calorimeters.

The momentum of cosmic ray muons traversing the entire ATLAS detector are measured twice, first in the top part of the muon spectrometer, then in the bottom part of the muon spectrometer. The comparison of the two momentum measurements allowed us to measure the fractional momentum resolution of the muon spectrometer up to muon with $p_T = 300 \text{ GeV}/c$ (see Figure 1.7). The measured momentum resolution is in agreement with the expected momentum resolution for $p_T \lesssim 100 \text{ GeV}/c$, but is a factor 2-3 worse than ex-

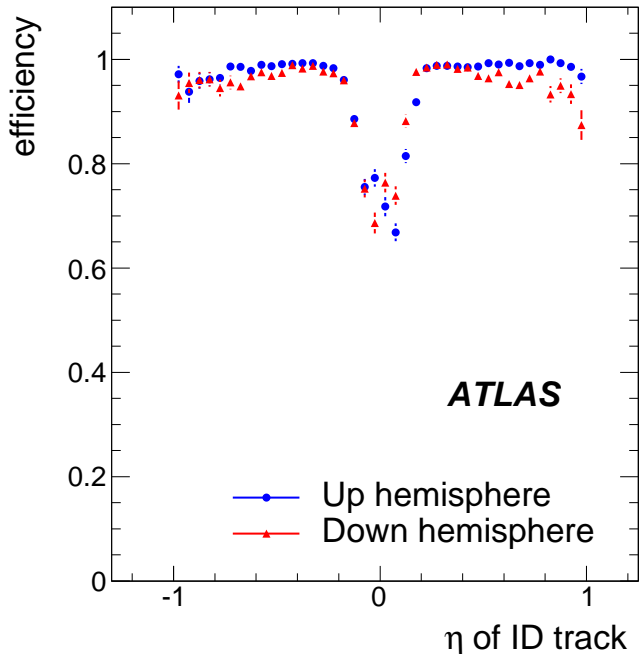


Figure 1.6: Track reconstruction efficiency in the muon spectrometer as a function of the pseudorapidity of the cosmic ray muon reconstructed by the inner detector. The muon momentum in the inner detector was requested to be greater than 5 GeV for the top part and greater than 9 GeV for the bottom part. The loss of efficiency in the region near $|\eta| = 0$ is due to acceptance holes of the muon spectrometer needed for services of the inner detector and the calorimeters.

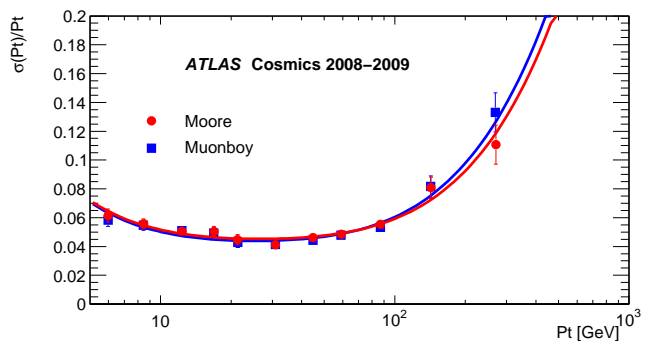


Figure 1.7: Transverse momentum resolution of the muon spectrometer evaluated by comparing the momentum of cosmic ray muons measured in the top part of the spectrometer with the momentum in the bottom part of the spectrometer in the barrel region of the muon spectrometer. The results of two complementary track reconstruction algorithms show similar performance.

pected above $300 \text{ GeV}/c$ due to the residual misalignment of the muon spectrometer.

The muon reconstruction efficiency and the momentum scale and resolution will be measured with $Z \rightarrow \mu^+ \mu^-$ decays in pp collision data at the LHC

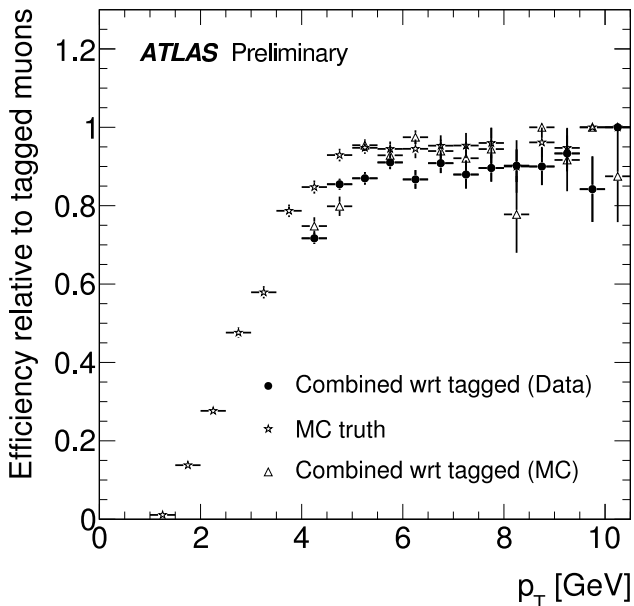


Figure 1.8: Comparison of the relative efficiency of the standard muon reconstruction algorithm with respect to reference muons as defined in the text with the Monte-Carlo prediction for $p_T > 4$ GeV/c. The star-shaped symbols labelled as MC truth show the standard muon reconstruction efficiency for muons identified as true muons in the simulation.

[30, 31]. The position of the Z resonance peak is a measure for the momentum scale, its width a measure for the momentum resolution. The muon reconstruction efficiency will be measured by the so-called "tag-and-probe method". In the tag-and-probe method events with an isolated muon and an inner detector track giving an invariant mass compatible with the Z boson mass are selected. The fact that one of the two tracks is identified as muon ensure that the second track is also a muon and one can count how often the second track is reconstructed as a muon. The method has been studied on Monte-Carlo data and is able to reproduce the muon reconstruction efficiency with a systematic uncertainty of 0.2%.

In the early phase of the LHC operation the rate of dimuon decays of Z bosons is too low for the tag-and-probe method and alternative methods have to be used to get a handle on the muon spectrometer reconstruction efficiency. The efficiency of the standard muon reconstruction is measured with respect to an alternative muon identification approach. In the alternative approach an inner detector track which deposits only little energy in the calorimeters and which can be matched with a track segment in the muon spectrometer is identified as muon. The relative efficiency of the

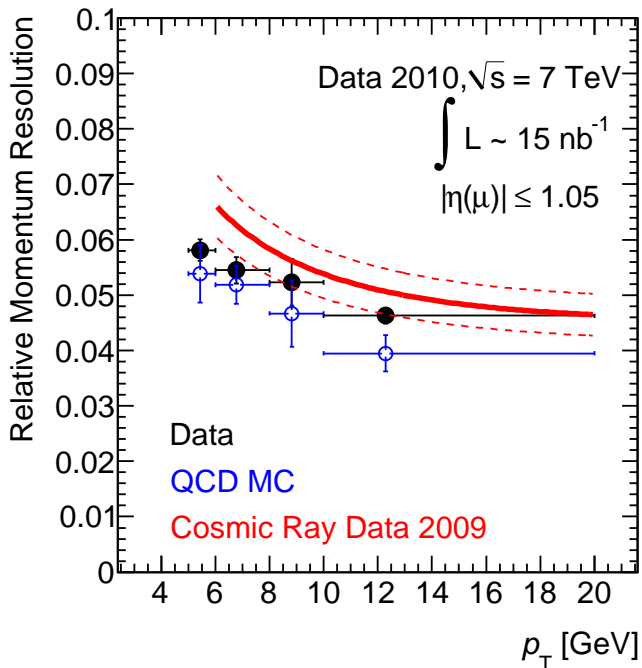


Figure 1.9: Comparison of the momentum resolution in the barrel part of the muon spectrometer measured with muons from pp collisions in comparison with the expected resolution from Monte-Carlo data and the analysis of cosmic ray data. The measured resolution agrees with the prediction within the measurement uncertainties.

standard muon reconstruction with respect to the alternative selection is shown as a function of the transverse muon momentum in Figure 1.8 [32]. The relative efficiency agrees with the Monte-Carlo prediction within errors. It is lower than the true efficiency as the alternative muon selection has a lower rejection power against muons from pion and kaon decays in flight than the standard reconstruction.

At transverse momenta below 20 GeV/c the momentum resolution in the muon spectrometer is significantly larger than in the inner detector and can be measured by comparing the momentum measured in the muon spectrometer with the momentum measured in the inner detector. The resolution measured this way is in agreement with the Monte-Carlo prediction and the measurement with cosmic ray muons within the measurement uncertainties (see Figure 1.9) [33].

The analysis of cosmic ray data and the first results of the analysis of pp collision data confirm the expected performance of the muon spectrometer.

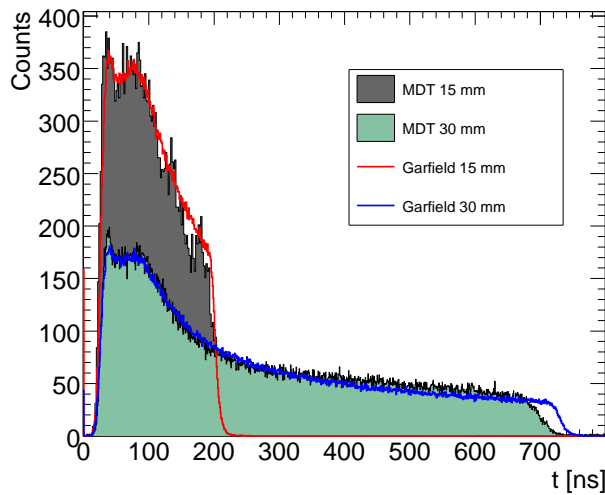


Figure 1.10: Comparison of the drift-time spectra of 30 mm and 15 mm diameter drift tubes with Ar:CO₂ (93:7) gas mixture at 3 bar and the same gas amplification of 20,000 which are in good agreement with the simulations with the GARFIELD program. The 15 mm diameter tubes show the same drift time spectrum as the 30 mm diameter tubes up to the cut off at the maximum drift time of 180 ns which corresponds to the smaller tube diameter.

1.1.3 Detector Development for ATLAS Upgrades

For the searches for the Higgs boson and other new phenomena at the LHC, as high as possible luminosity of the accelerator is desirable. The design luminosity of the LHC of $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ is expected to be reached after a few years of data taking at lower luminosity. With a series of upgrades of the CERN accelerator complex a further increase of the LHC luminosity by an order of magnitude can be envisaged (Super-LHC or SLHC). Then, at the latest, a major upgrade of the ATLAS detector will be necessary including parts of the muon spectrometer, in particular in the forward regions with respect to the proton beams where the background rates will be the highest. The planning of a possible detector upgrade has started already within the ATLAS collaboration. In this context, we have developed a research program and submitted proposals for the development of improved muon drift tube detectors [34], of radiation tolerant readout electronics [35] and of methods to reduce the data volume [36] or increase the bandwidth of the data acquisition system of the MDT chambers [35] for operation of the ATLAS detector at very high luminosities.

Already at the design luminosity, the MDT chambers have to cope with unprecedentedly high back-



Figure 1.11: Front view of bundles of drift tubes with 30 mm (left) and 15 mm (right) diameter.

ground count rates from neutrons and photons as byproducts of the proton-proton collisions. The effects of high count rates in the MDT chambers of up to 500 Hz/cm^2 at the design luminosity in the inner layers of the forward regions have been studied by our group in a muon beam at the Gamma Irradiation Facility at CERN [13]. In addition to a degradation of the spatial resolution of the drift tubes from $80 \mu\text{m}$ without background to about $100 \mu\text{m}$ at the highest rates caused by fluctuations of the increasing space charge in the tubes, the muon detection efficiency of the tubes deteriorates quickly from 94% to 70%. New muon chambers with faster response time and high track reconstruction efficiency are required for the ten times higher background rates expected at a Super-LHC in the spectrometer endcap regions, and are desirable for the highest-rate regions already at the LHC design luminosity.

The development of faster muon detectors for Super-LHC at MPI builds on the extensive experience with the design and operation of the present ATLAS MDT chambers. First tests [37] confirmed the expectation that by reducing the drift tube diameter by a factor of two while leaving the gas mixture, pressure and gain unchanged, the maximum drift time is reduced by a factor of 3.5 (see Fig. 1.10). At the same time, the count rates caused by neutron and photon conversions in the drift tube walls will decrease proportional to the tube circumference by a factor of two, reducing the drift tube occupancy (the product of maximum drift time and count rate) and, consequently, the efficiency

loss by a factor of 7. In addition, the number of drift tube layers fitting into the same volume as the present muon chambers can be increased by up to a factor of two (see Fig. 1.11) improving the chamber resolution and track reconstruction efficiency.

A concept for the drift tube and chamber design has been developed with the aim to construct and test a prototype chamber in 2008. The development of new radiation hard readout electronics for the MDT chambers has also started.

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1.2 ATLAS Physics Analysis

Parallel to the construction of the detector components the MPI ATLAS group started the preparation of physics analysis of hadron collision data at the LHC. The results obtained in the years 1997-2006, including preparatory work based on Tevatron data, are described in the previous reports [1, 2], and references therein.

The present physics studies for the ATLAS experiment cover a broad physics range. Already at an early stage of data taking, a number of Standard Model processes will occur in abundance. These processes allow for detailed studies of the detector performance, as well as the precision measurement of QCD and electroweak observables, the top-quark physics being of particular interest. Good understanding of the Standard Model processes is essential also for the new discoveries. The ATLAS discovery potential is explored in the searches for the Higgs boson both in the Standard Model and in supersymmetric extensions, as well as in a generic search for supersymmetric particles.

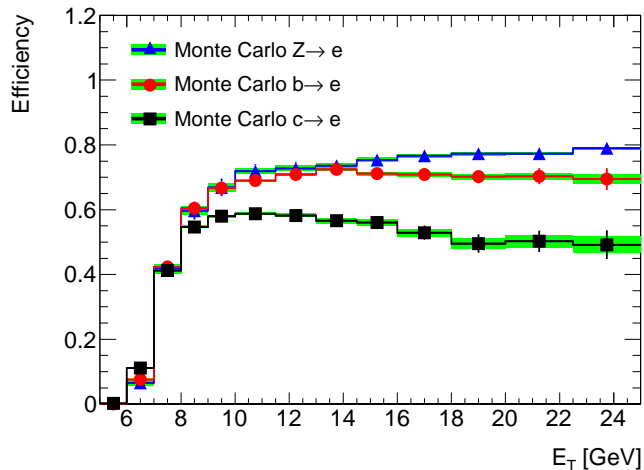


Figure 1.12: Monte-Carlo prediction of electron reconstruction efficiency for electrons from heavy quark and Z boson decays after the optimization of the electron identification cuts.

The ongoing studies are described in more detail below.

1.2.1 Measurement of inclusive lepton cross sections

At the LHC pp collision events with highly energy electrons and muons in the final provide clean signatures for many physics processes of interest. A good understanding of the inclusive electron and muon cross section is therefore of great importance. The MPI group contributes to the measurement of the inclusive electron and muon cross section.

At the LHC electrons are produced predominantly produced in decays of heavy quarks for transverse momenta below about 30 GeV/c and in decays of W and Z boson at higher transverse momenta. The MPI group significantly contributed to the optimization of the electron selection criteria to arrive at an electron selection efficiency which is flat in the transverse electron momentum (see Figure 1.12). The first measured inclusive electrons spectrum at a centre-of-mass energy of 7 TeV at the LHC shown in Figure 1.13 agrees with the prediction of the Pythia minimum bias Monte-Carlo prediction within about 20%. The MPI group contributes to the analysis of this discrepancy with increasing statistics of the inclusive electron sample.

The MPI group is also involved in the measurement of the inclusive muon cross section contributing with its experience in muon performance studies. The measured inclusive muon p_T spectrum is presented

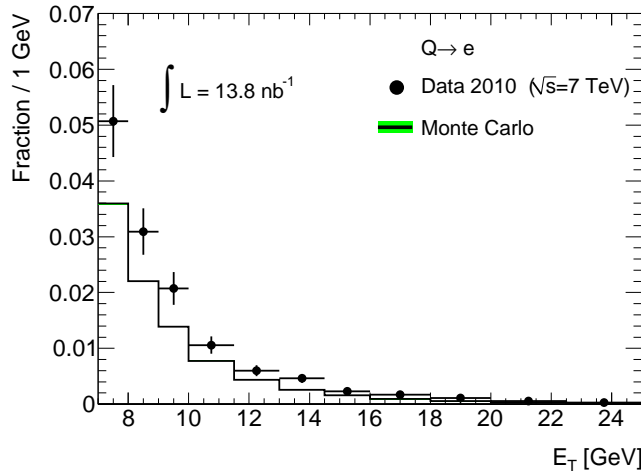


Figure 1.13: Comparison of the measured distribution of the transverse energies of prompt electrons from c and b decays with the Pythia 6.4 minimum bias Monte-Carlo predictions.

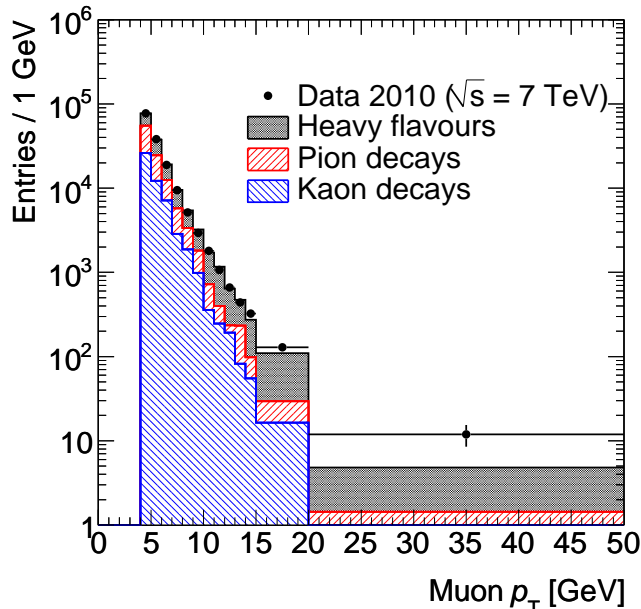


Figure 1.14: Comparison of the measured inclusive muon transverse momentum spectrum with the Pythia 6.4 minimum bias Monte-Carlo prediction. The Monte-Carlo data is decomposed into the three sources of muons, namely in-flight decays of charged pions and kaons and the decays of heavy-flavour hadrons.

in Figure 1.14 where it is compared with the Pythia 6.4 minimum bias Monte-Carlo prediction. The measured spectrum is well reproduced by the Monte-Carlo simulation. The main sources of the muon at transverse momenta below 20 GeV/c are in-flight decays of charged pions and kaons and the decays of heavy-flavour hadrons according to the Monte-Carlo simulation. Decays of W and Z bosons become important

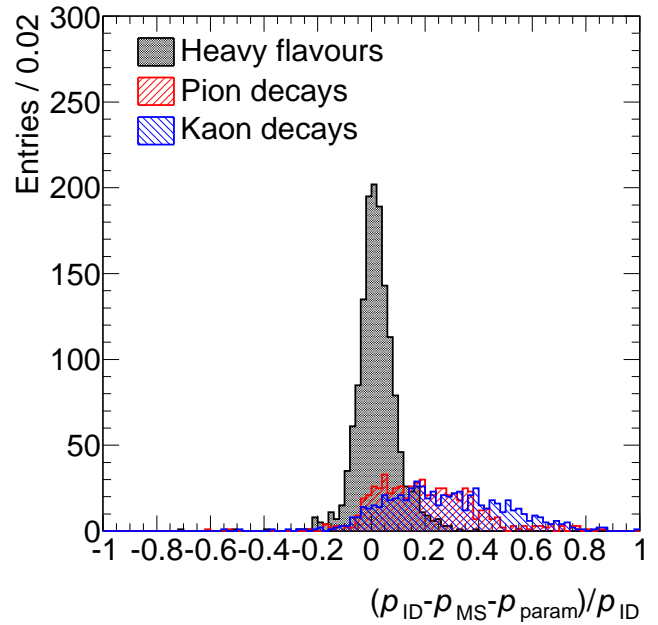


Figure 1.15: Distribution of the difference of the muon momentum measurements in the inner detector and the muon spectrometer normalized to the inner detector momentum measurements in simulated data. Decays of pions and kaons in flight lead to a large tail to positive values as the inner detector measures the pion and kaon momentum while the muon spectrometer measures the momentum of the decay muon.

for $p_T > 20$ GeV/c. The contribution of pion and kaon decays in-flight to the inclusive muon spectrometer will be estimated from data by comparing the momentum measured in the inner detector with the momentum measured in the muon spectrometer. For late pion and kaon decays in the inner detector lead to a large momentum imbalance between the inner detector and muon spectrometer as illustrated in Figure 1.15.

1.2.2 Study of Standard Model Processes

Standard Model processes are studied with the motivation of estimating the backgrounds to the production of Higgs bosons and supersymmetric particles from the data and to improve the Monte Carlo simulations. The processes under investigation are QCD high p_T and multi-jet production, forward jet production, b jet production, inclusive W , Z and $t\bar{t}$ production as well as $Z \rightarrow \tau^+\tau^-$ decays. $Z \rightarrow e^+e^-$ and $\mu^+\mu^-$ decays are being studied for the purpose of detector calibration and data quality monitoring. The analysis methods developed will be applied to first measurements of multi-jet, $b\bar{b}$ and inclusive W and Z production cross-sections with the early data.

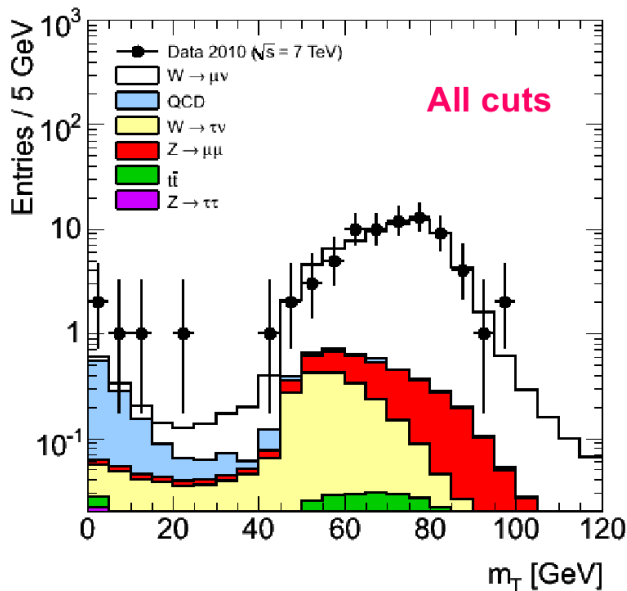


Figure 1.16: Transverse mass distribution of $W \rightarrow \mu\nu_\mu$ candidates found in the first 16 nb^{-1} of pp collision data collected by ATLAS.

Figure 1.16 shows the transverse mass distribution of the $W \rightarrow \mu\nu_\mu$ candidates in the first 16 nb^{-1} of pp collision data collected by ATLAS. The measured distribution is consistent with the Monte-Carlo prediction and has a negligible background contamination. In the same set of pp collision data ATLAS has observed 8 Z boson candidates as can be seen as an excess of entries at the Z mass in the dimuon mass spectrum of Figure 1.17.

1.2.3 Top-Quark Physics

The top-quark is by far the heaviest elementary building block of matter. Therefore it should have the strongest couplings to any mechanism that generates mass, which makes it a very interesting object for an unbiased search for this mechanism. In addition, the precise knowledge of the quantum numbers of the top-quark helps to further constrain the parameters of the Standard Model, and is a mandatory prerequisite for any study of new physics that will almost inevitably suffer from top-quark reactions as background processes.

The main interest of the top-physics analyses work of the MPI group is the investigation of the $t\bar{t}$ production process, and particularly the determination of the mass of the top-quark and the $t\bar{t}$ production cross-section in the reaction $t\bar{t} \rightarrow b\bar{b}W^+W^-$. The emphasis in these two measurements is rather different. For

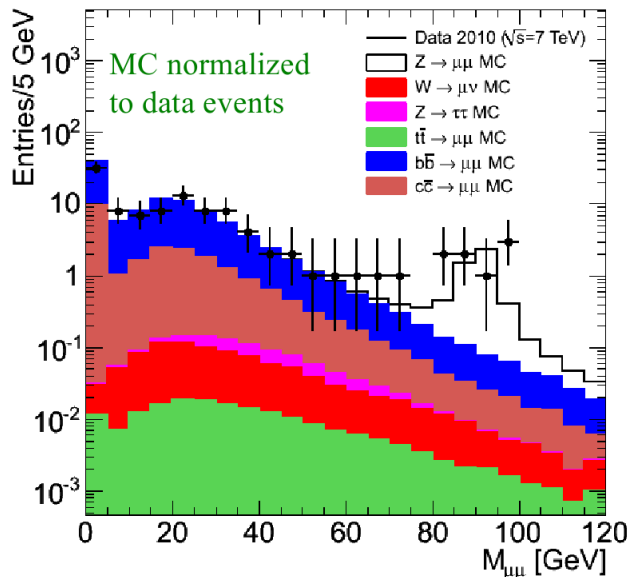


Figure 1.17: Invariant dimuon invariant mass distribution for isolated muons in the first 16 nb^{-1} of pp collision data collected by ATLAS. An excess of 8 Z events is visible in the distribution consistent with the Monte-Carlo prediction.

the mass measurement all systematic effects that may change the obtained mass and resolution are important, in contrast, for the cross-section measurement the main emphasis is on the detailed understanding of all efficiencies.

The analyses use two decay channels of the W -boson pair, the semi-leptonic channel, where the W -boson pair decays into $\ell\nu qq'$ with $\ell = e, \mu$, (30% branching ratio) and the full-hadronic channel, where both W -bosons decay into a qq' pair (44% branching ratio). In both channels, the mass of the top-quark will be obtained from hadronically decaying W -bosons and the corresponding b -jet.

The main background reactions, as determined from Monte Carlo simulations, are the $W+N$ -jets production and that fraction of the $t\bar{t}$ production, where the W -boson pair decays into the other decay channels. An additional potential background process is the QCD multi-jet production, especially for the analyses of the full-hadronic channel. For this background, in the semi-leptonic channel, the lepton candidate stems either from a semi-leptonic decay of a heavy quark of the QCD multi-jet final state, or from instrumental background. Before any selection, this background is huge, such that event samples covering the full phase space cannot be simulated, and eventually this background contribution has to be obtained from the data themselves. At present, only initial studies of this back-

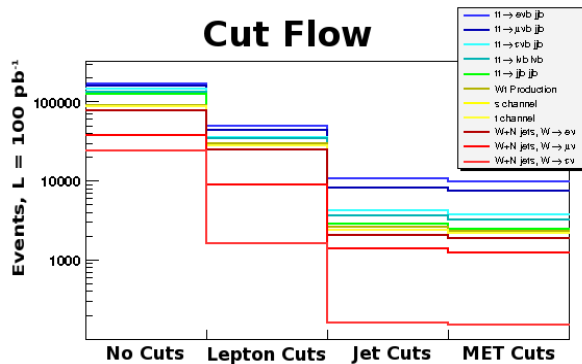


Figure 1.18: Event composition in the semi-leptonic sample.

ground have been performed, and it is not included in the numbers quoted below. In the semi-leptonic channel the charged leptons with high transverse momenta from the decay of one W-boson are utilized to identify the event and to efficiently suppress background from the QCD multi-jet production. In the full-hadronic channel only kinematic requirements can be used, therefore, this channel suffers from a much higher background from QCD multi-jet production.

The event selection for the semi-leptonic channel requires an isolated electron or muon within the good acceptance of the detectors, which has a transverse momentum of more than 20 GeV and lies within the rapidity range of $|\eta| < 2.5$. To account for the neutrino, a missing transverse energy of more than 20 GeV is required. In addition, at least four jets are required within the same range of rapidity, and having transverse momenta of more than 40 GeV for the three highest p_T -jets, and more than 20 GeV for the fourth jet. All jets should be well separated from the identified lepton. The event composition for the main steps of the selection are shown in Fig. 1.18. With this selection, an average signal efficiency of about 20 (30)% is reached for the electron (muon) samples, and the signal to background ratio improves by a factor of ten to about 1.6. For the full-hadronic channel the requirements on the isolated lepton and the missing transverse energy are inverted. In addition, at least six jets, well separated from electrons, are required within the same region of rapidity, and having transverse momenta of more than 20 GeV. Here, the selection efficiency is about 25% but the amount of QCD and W+N-jets background has yet to be evaluated.

The $t\bar{t}$ production is not only of prime physics interest, but also presents a very clean event sample to improve on the reconstruction algorithms. Of key im-

portance are the b-tagging performance and the calibration of the jet energy scale, which at present constitutes the most important uncertainty on the mass of the top-quark as determined at the Tevatron collider. Since a rather clean sample of $t\bar{t}$ events can be selected without employing b-tagging, and the b-quarks are mostly the highest p_T jets, the performance of the b-tagging algorithms can be studied in a rather unbiased sample. The jet calibration can make use of the fact that the invariant mass of the qq' pair stemming from the W-boson decay has to match the known W-boson mass.

The steering of the b-tagging algorithm is a compromise of efficiency and purity, while attaining an acceptable rejection rate for jets originating from light quarks. This light jet rejection rate is most important for suppressing events from QCD multi-jet production, which is dominated by jets from light quarks. Similarly, it helps to suppress the W+N-Jets background. The b-tagging performance has been studied for a standard ATLAS algorithm as a function of the jet algorithms, when requesting a 50% b-tagging efficiency for all collections. In this situation, the light jet rejection varies from about 100 to 700 with the inclusive k_T -jet algorithm for $R = 0.4$ (see below for a definition) giving the best performance.

Both semi-leptonic and full hadronic top-pair events allow for a so-called in-situ calibration of the jet energy scale. The known mass of the W can be used to constrain the energy of the light quark jets from the hadronic W decays either to apply a final calibration step especially in early data taking where the uncertainty in the jet energy scale is expected to be large, or to verify the applied calibration.

Based on generated $t\bar{t}$ events from the MC@NLO Monte Carlo program, and W+N-jets as predicted by the Alpgen generator, an analysis has been performed to obtain the mass of the top-quark in the semi-leptonic channel. For each event, the four jets with the highest transverse momenta are assumed to stem from the $b\bar{b}$ and qq' systems. They are grouped to obtain the estimate of the masses of the hadronically decaying W-boson and the associated top-quark. At present, the three jet combination which maximizes the transverse momentum is chosen to form the hadronically decaying top-quark. Out of this, the two jet combination that results in an invariant mass closest to the known W-boson mass is taken to represent the W-boson. In future other choices will be compared to this. The spectrum in Fig. 1.19 is fitted with a Gaussian function to parameterize the correct combinations leading to the

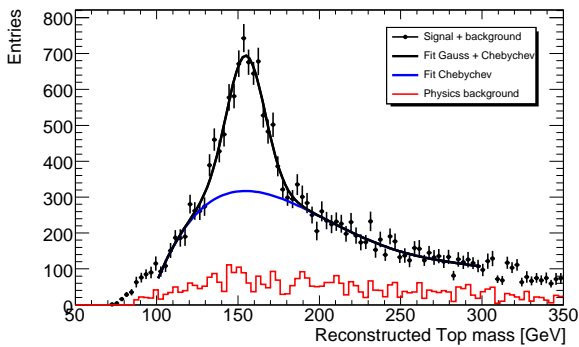


Figure 1.19: Fit to reconstruct the top-quark mass for a sample containing $t\bar{t}$ and W+N-Jets events.

top-quark mass and width, and a sum of Chebyshev polynomials used to describe the background stemming from the W+N-jets events and wrong jet combinations in signal events. This figure shows the result when using the inclusive k_t -jet algorithm and $R = 0.4$. The central value of the fitted spectrum as well as the width of the distribution strongly depend on the choice of the jet algorithm and jet steering parameters, the calibration of the calorimeters, and various other details of the analysis. Using the present status of the knowledge of the jet energy scale, a detailed study of the dependence of the reconstructed top-quark mass on the jet algorithms and their parameters has been performed [4] using only $t\bar{t}$ events where at least on W-boson decays leptonically, and with a top-quark mass of 175 GeV. Fig. 1.20 shows the results for the cone-jet algorithm, and the k_t -jet algorithm in the E recombination scheme, and using both the inclusive and the exclusive mode, as function of the parameters of the jet algorithms. The abbreviations at the horizontal axis denote the algorithm and the values for the steering parameters, i.e. the value of R for the radius of the cone-algorithm and the distance parameter R for the k_t -jet algorithm. The exclusive k_t -jet algorithm modes are indicated by the invariant mass cut, e.g. $D = 5^2 \text{ GeV}^2$, or the requested jet multiplicity, e.g. $N = 4$. All collections use calibrated top cluster as input, the * denotes an additional calibration based on cone-jets with $R = 0.4$. In this study, for the first time the k_t -jet algorithm has been used in top physics analyses at ATLAS. Because of its better stability against divergences, this algorithm is theoretically preferred over the traditionally applied cone-jet algorithm. One result of the investigation is that values of the R -parameter of $R = 0.4 - 0.5$ (Kt4 and Kt5 in Fig. 1.20) are much preferred over the default

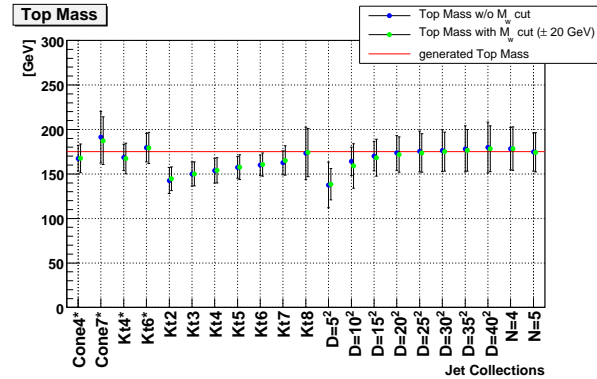


Figure 1.20: Top-quark mass value and uncertainty as a function of the parameters of the jet algorithms.

value of $R = 1$ (not shown anymore), since for the latter choice for many events the two jets stemming from the W-Boson decay are merged into one reconstructed jet, thereby resulting in a loss of the event. On the vertical axis Fig. 1.20 the mean value and the sigma of the Gaussian are shown for the standard selection described above, and when using the additional requirement that the reconstructed mass of the W-boson coincides with the known mass to within 20 GeV. A good performance is signalled by small variations of the mean value when changing the steering parameter, and a small ratio of width and mean. In contrast, the absolute value mainly reflects the present level of calibration, which can be seen e.g. be comparing Kt4 and Kt4*. This means the best performance is achieved with the inclusive k_t -jet algorithm and medium values for R , around $R = 0.4 - 0.5$. The requirement on the W-boson mass does not improve on the mass determination. Clearly, the final value of the R -parameter needs to be found once the calibration of the hadronic calorimeters is finalized.

As stated above, the jet assignment to the W-boson and top decay products is not unique and also not very efficient. For comparison, at the Tevatron it has been found that the correct jet-to-quark assignment is only achieved in about 40% of the cases. This led to the development of the matrix element method in which, rather than choosing one particular combination, each combination is considered, and weighted by the corresponding probability, given by the matrix element. In order to better understand the jet assignment and resolution, and to find new observables to further suppress the W+N-Jets background, an analysis has been performed investigating the correlation of jets reconstructed from stable hadrons to jets reconstructed from calo topo clusters. Here, a pair of jets is called matched

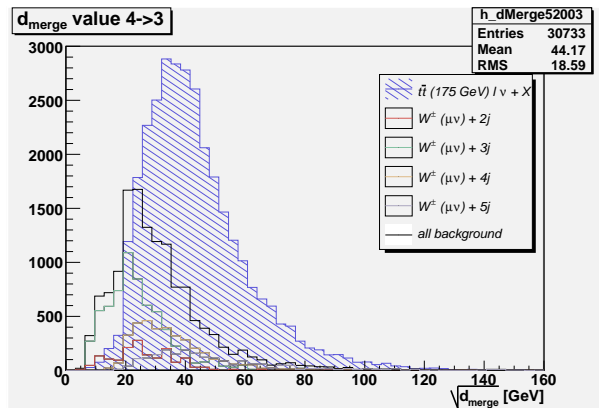


Figure 1.21: The $d_{\text{merge}}(4 \rightarrow 3)$ distribution for signal and background.

if for a stable hadron jet the closest reco jet falls within a distance of $\Delta R = 0.15$, where each reco jet can only be assigned to one stable hadron jet. For the signal sample, it has been found that the overall matching efficiency is 92%, rather constant in rapidity but with a steep turn-on at low transverse momenta.

The separation of jets can be monitored when running the k_T -jet algorithm by studying the $d_{\text{merge}}(M \rightarrow M - 1)$ values at which an M-Jet configuration is reduced to an (M-1)-Jet configuration. For sufficiently different jet structure in signal and background events, this event shape variable can be used as discriminator. As an example, the $d_{\text{merge}}(4 \rightarrow 3)$ distribution is shown in Fig. 1.21 for signal and W+N-Jets background events with $N = 2 - 5$ and $W \rightarrow \mu\nu$ decays. For this distribution, the semi-leptonic event selection has been performed. Clearly, the tail at low values is more pronounced for the background events, e.g. requiring $d_{\text{merge}}(4 \rightarrow 3) > 20 \text{ GeV}$ reduces the signal efficiency by only about 2%, while improving the purity by about 5%. The implication on the reconstructed top-quark mass is under study.

Concerning the cross-section measurement a self consistent analysis has been performed for the semi-leptonic channel by declaring one tenth of the signal and background samples as data, and using the remaining events to obtain the signal efficiency and perform the background subtraction. With this, the input cross-section could be recovered and the analysis now moves on to obtain systematic uncertainties.

Inevitably, one main initial uncertainty will come from the luminosity determination, which is provided by external measurements. Another important ingredient is the trigger efficiency. For the semi-leptonic channel it is found that an efficient trigger is provided

by the leptons of high transverse momenta. The trigger efficiency for the signal sample is about 80 (90)% for the electron (muon) sample, rather flat in rapidity (when excluding crack regions), and with a steep turn-on at transverse momenta around 20 GeV. A subsample of these events is also triggered by jet triggers, allowing for an unbiased determination of the lepton trigger efficiencies. Utilising this, the next step is to obtain a trigger efficiency grid in transverse momenta and rapidity of the leptons for correcting the data.

In conclusion, initial analyses have been performed for both the mass and the cross-section measurements investigating some sources of systematic uncertainties. The future work will concentrate on investigations of other sources of systematic uncertainties.

1.2.4 Searches for the Higgs Boson

The Standard Model Higgs Boson

The origin of particle masses is one of the most important open questions in particle physics. In the Standard Model, the answer to this question is connected with the prediction of a new elementary particle with spin 0, the Higgs boson H. The discovery of the Higgs boson, the last missing particle of the Standard Model, is the main motivation for the LHC experiments. The mass of the Higgs boson is not predicted by the theory. The present experimental lower bound of $115 \text{ GeV}/c^2$ and the theoretical upper limit of about $1 \text{ TeV}/c^2$ require a wide mass range to be explored. The preferred Standard Model Higgs mass range derived from the precision measurements of the parameters of the electroweak gauge theory is below $200 \text{ GeV}/c^2$.

In the mass range above $180 \text{ GeV}/c^2$, the key discovery channel which also provides a direct measurement of the Higgs boson mass is the Higgs decay into four charged leptons via two intermediate Z bosons. The lower mass range can be covered by searches for several Higgs decay modes (see Fig. 1.22). Although the ATLAS detector has been designed for a discovery of the Higgs boson over the full mass range, Higgs boson searches at the Large Hadron Collider remain challenging because of the high cross sections of the background processes exceeding the signal by many orders of magnitude. Selective triggers, efficient background suppression and reliable prediction of the background contributions are essential.

Our present studies focus on the preparation for an early discovery of the Higgs boson during the first years of running at the initial LHC luminosity

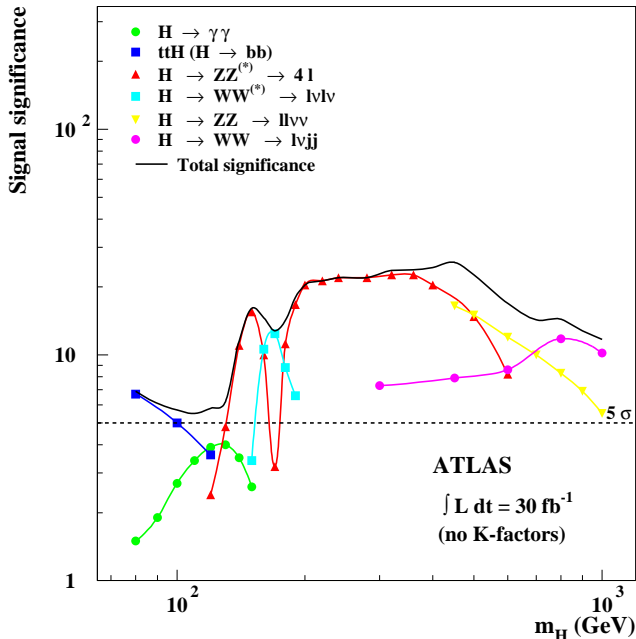


Figure 1.22: Discovery potential of the ATLAS experiment for the Standard Model Higgs boson. The statistical significance expected for an integrated luminosity of 30 fb^{-1} is shown for the different Higgs decay modes discussed in the text and their combination as a function of the Higgs boson mass in the range below $200 \text{ GeV}/c^2$ preferred by previous precision tests of the Standard Model. The grey shaded region has been excluded by direct Higgs boson searches of the experiments at the LEP e^+e^- collider.

of $10^{33} \text{ cm}^{-2}\text{s}^{-1}$. In the first year of LHC operation, accumulation of a data sample corresponding to an integrated luminosity of $1 - 10 \text{ fb}^{-1}$ is expected. Within three years of stable running of the experiment, 30 fb^{-1} of data can be collected. As described in the following, most of the Higgs boson decay channels which are promising for a discovery are investigated by our group. Detection and comparison of signals in several decay modes will be crucial to confirm the existence of the Higgs boson and to distinguish between the predictions of the Standard Model and of its possible extensions. The analyses are performed with detailed Monte Carlo simulations of the signal and background processes in the ATLAS detector using the latest description of the detector geometry and performance.

In the mass region below $120 \text{ GeV}/c^2$, the Higgs boson predominantly decays into $b\bar{b}$ pairs. Due to large QCD background in the gluon-fusion production mode, this decay can only be triggered and discriminated from the background in the production mode of the Higgs boson in association with a $t\bar{t}$ pair. Our

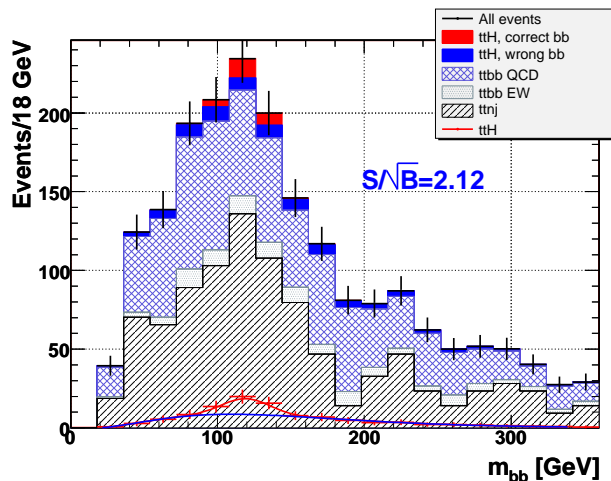


Figure 1.23: Invariant mass distribution of reconstructed $b\bar{b}$ jet pairs in simulated $pp \rightarrow t\bar{t}H, H \rightarrow b\bar{b}$ events with a Higgs mass of $120 \text{ GeV}/c^2$ together with the contributions from $t\bar{t}b\bar{b}$ and $t\bar{t}jj$ background for an integrated luminosity of 30 fb^{-1} . The red and the blue curves show the correct and the wrong $b\bar{b}$ combinations in the signal events, respectively.

analysis [5] uses a neural network based discrimination method to study the observability of the $t\bar{t}H \rightarrow (\ell\nu b)(q\bar{q}b)(b\bar{b})$ channel, where one top quark decays semileptonically providing a high momentum lepton trigger and the other one hadronically. Four bottom quark jets have to be identified in these complex final states.

Fig. 1.23 shows the invariant-mass distributions of the correct and the wrong $b\bar{b}$ combinations in the signal events, as well as the contributions from irreducible ($pp \rightarrow t\bar{t}b\bar{b} + X$) and reducible ($pp \rightarrow t\bar{t}jj + X$) background. The signal significance $S/\sqrt{B} \approx 2$ determined from the numbers of signal (S) and background (B) events in a $2\sigma_m$ mass window around the simulated Higgs mass of $120 \text{ GeV}/c^2$ does not include systematic uncertainties. Due to the relatively low efficiency of the b jet reconstruction and the similarity of the $b\bar{b}$ invariant mass distributions for signal and background, high integrated luminosity and a precise determination of the Standard Model background contributions from data is needed for an observation of the Higgs boson in this channel. We work on improvements of the signal discrimination from the background and on methods to estimate the background contributions from the first LHC data which are essential for the detection of the Higgs boson in this decay channel.

A feasibility study has been performed of the detection of $H \rightarrow b\bar{b}$ decays in the Higgs production process by W or Z gauge boson fusion [6]. As for other Higgs decay modes (see below), the associated forward-jet production and corresponding central jet suppression provide signatures for additional background rejection. The trigger efficiency for this process with standard ATLAS lepton and jet triggers was found to be acceptable. However, the background rate especially from direct $b\bar{b}$ production, is too high for identifying a signal in this channel.

The second most frequent Higgs decay mode which can be observed in the mass range below $140 \text{ GeV}/c^2$ is the decay into a $\tau^+\tau^-$ pair. This decay can only be discriminated against background processes in the Higgs production mode via vector-boson fusion where two additional forward jets in the final state provide signature for background rejection. The decay mode with both τ leptons decaying leptonically as well as with a hadronic and a leptonic τ decay have been studied. The selection criteria have been optimized using multivariate analysis techniques. The $\tau^+\tau^-$ invariant mass is determined under the approximate assumption that all τ decay products are emitted in the direction of the τ lepton. The expected invariant mass distribution for the signal of a $120 \text{ GeV}/c^2$ Higgs boson against backgrounds is shown in (Fig. 1.24).

The mass range above $130 \text{ GeV}/c^2$ can be covered with the decay $H \rightarrow W^+W^- \rightarrow (\ell\nu)(\ell\nu)$ both in gauge-boson fusion and, at higher masses, in gluon fusion production. Parallel to the optimization of the event selection for these channels, we are studying the performance of the forward-jet reconstruction and of the veto against central jet production, which are used for the suppression of the $t\bar{t}$ and WW +jets backgrounds to the vector boson fusion production processes. A new jet reconstruction algorithm using particle tracks in the inner detector instead of energy deposition in the calorimeters has been developed [7]. The tracks forming a jet can be associated to a common vertex suppressing the reconstruction of spurious jets in the presence of multiple proton-proton interactions per beam collision (pile-up events) and, thereby, improving the efficiency of the central jet veto.

While this decay channel allows for fast Higgs boson discovery with an integrated luminosity of less than 10 fb^{-1} , it cannot provide a precise measurement of the Higgs mass because of the two neutrinos in the final state. Precise determination of the background contributions is, therefore, required to observe an ex-

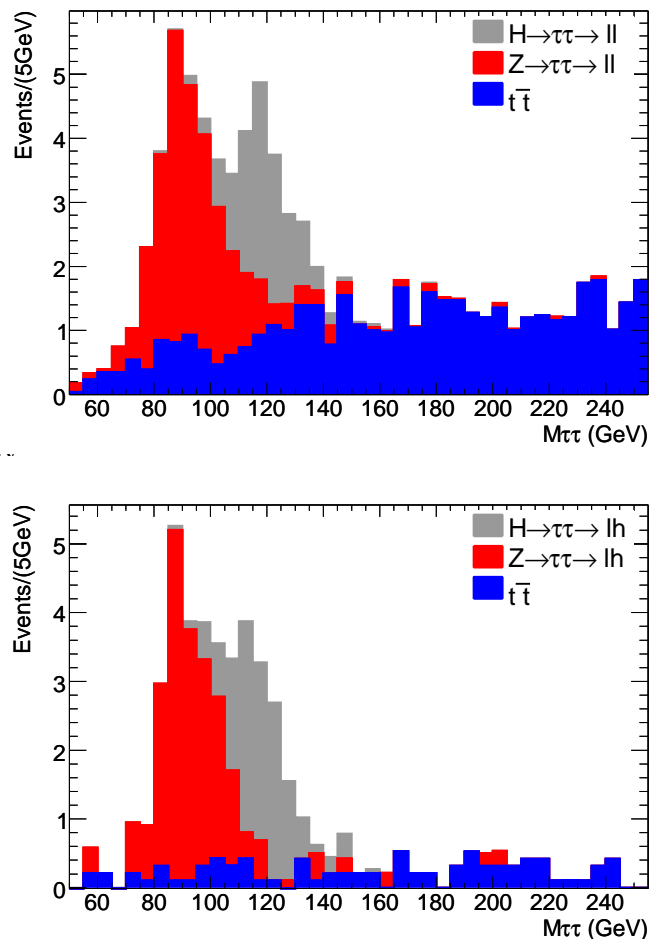


Figure 1.24: Invariant τ pair mass distribution in simulated $H \rightarrow \tau^+\tau^-$ decays in vector-boson fusion production with a Higgs mass of $120 \text{ GeV}/c^2$ together with background contributions from $t\bar{t}$ and Z +jets production for an integrated luminosity of 30 fb^{-1} . The two τ leptons have been reconstructed either both in purely leptonic decays (top picture) or one in leptonic and the other one in hadronic decay mode (bottom picture).

cess of signal events. For this purpose, we are studying the measurement of the Standard Model background processes with the early LHC data.

Above a Higgs mass of $130 \text{ GeV}/c^2$ the clearest signature is found in the four-lepton decay $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ which also allows for precise Higgs mass measurement. The reconstruction of this channel strongly relies on the high lepton-identification efficiency and good momentum resolution of the ATLAS detector. Our initial analysis of the four-muon final state [8] has been extended to include also the four-electron and $2e2\mu$ final states. The reducible $pp \rightarrow t\bar{t} + X$ background can be rejected by requiring lepton isolation and cutting on the masses of the reconstructed Z

bosons.

After the optimization of the cuts, the contribution of the $t\bar{t}$ and $Z + b\bar{b}$ backgrounds become small compared to the irreducible $pp \rightarrow ZZ^{(*)} + X$ background. Fig. 1.25 shows the four-lepton invariant mass distributions for $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ decays with 130 GeV/c and 300 GeV/c² Higgs mass and for the background processes corresponding to an integrated luminosity of 10 fb⁻¹ [9]. The first one is the most difficult case: the lower end of the sensitive mass range where the decay branching ratio becomes very small. For comparison, for a Higgs mass of 300 GeV/c², the $H \rightarrow Z$ decay reaches the maximum of its branching ratio and is the best channel for discovering the Higgs boson. Our recent studies include the precise determination of the background contributions from the data and the influence of pile-up events and radiation background in the ATLAS cavern.

Higgs Bosons Beyond the Standard Model

The minimal supersymmetric extension of the Standard Model (MSSM) predicts the existence of five Higgs bosons, a relatively light neutral one h behaving like the Standard model Higgs boson, two heavy neutral ones H and A and two heavy charged Higgs bosons H^\pm . Their production cross sections and decay branching ratios are determined by two independent parameters, e.g. the ratio $\tan\beta$ of the vacuum expectation values of the two Higgs doublets in this model and the pseudoscalar Higgs boson mass m_A where the production cross sections increase with $\tan\beta$ and fall with increasing m_A .

The neutral Higgs decay modes into two intermediate gauge bosons are suppressed in this model. On the other hand, the A or H boson decays into charged lepton pairs, $\mu^+\mu^-$ and $\tau^+\tau^-$, are enhanced compared to the Standard Model. The latter decay channel has an about three hundred times higher branching ratio than the first one but is more difficult to reconstruct and provides a less precise Higgs mass determination. Final states with both τ leptons with one hadronic and one leptonic τ decay have been studied.

In spite of the enhanced MSSM production rate, the signal events are still hidden under large $Z \rightarrow \mu^+\mu^-/\tau^+\tau^-$ and $t\bar{t}$ backgrounds. The detection of the b quark jets in the $gg \rightarrow b\bar{b}(A/H) + X$ production process, which dominates for large $\tan\beta$ and m_A , allows for strong background rejection. Additional background suppression is achieved by requiring large angular separation between the two decay

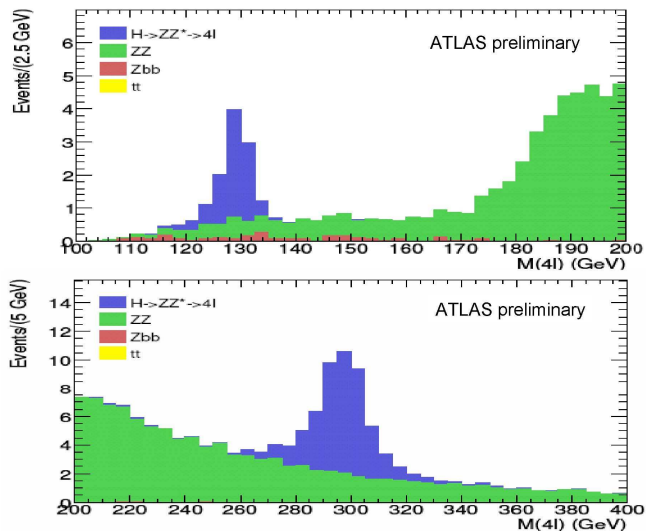


Figure 1.25: Four lepton invariant mass distribution in simulated $pp \rightarrow H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ events with a Higgs mass of 130 GeV/c² (top picture) and of 300 GeV/c² (bottom picture) together with the background contributions from $t\bar{t}$ and $ZZ^{(*)}$ production for an integrated luminosity of 10 fb⁻¹ [9].

leptons. Fig. 1.26 shows the invariant mass distributions of the lepton pairs in MSSM $A/H \rightarrow \mu^+\mu^-$ [10] and $A/H \rightarrow \tau^+\tau^-$ [11] decays together with the contributions from competing background processes corresponding to 30 fb⁻¹ of integrated luminosity. In Fig. 1.27, the regions of the $\tan\beta$ - m_A parameter space in which the $A/H/h \rightarrow \mu^+\mu^-$ and the $A/H \rightarrow \tau^+\tau^-$ signal, respectively, can be detected with a significance of at least 5 standard deviations are shown for an integrated luminosity of 10 fb⁻¹ [10],[11]. The discovery region for the $A/H \rightarrow \tau^+\tau^-$ decays is extended compared to the $A/H \rightarrow \mu^+\mu^-$ channel to smaller $\tan\beta$ values by about a factor of 5 at the lower end of the A boson mass range and more with increasing m_A .

The light neutral MSSM Higgs boson is difficult to distinguish from the Standard Model Higgs boson. Clear evidence for physics beyond the Standard Model would be provided by the discovery of charged scalar Higgs bosons. We have studied the search for the charged MSSM Higgs bosons in the decay channel $H^\pm \rightarrow \tau^\pm\nu_\tau$ which dominates for relatively small Higgs masses below 200 GeV/c². The charged Higgs bosons are produced in top quark decays in $pp \rightarrow t\bar{t} \rightarrow (bH^\pm)(bW^\mp)$ events. The τ leptons from H^\pm decays are reconstructed in their hadronic decay modes while the W bosons from top quark decays are required to decay leptonically. Since the H^\pm mass cannot be re-

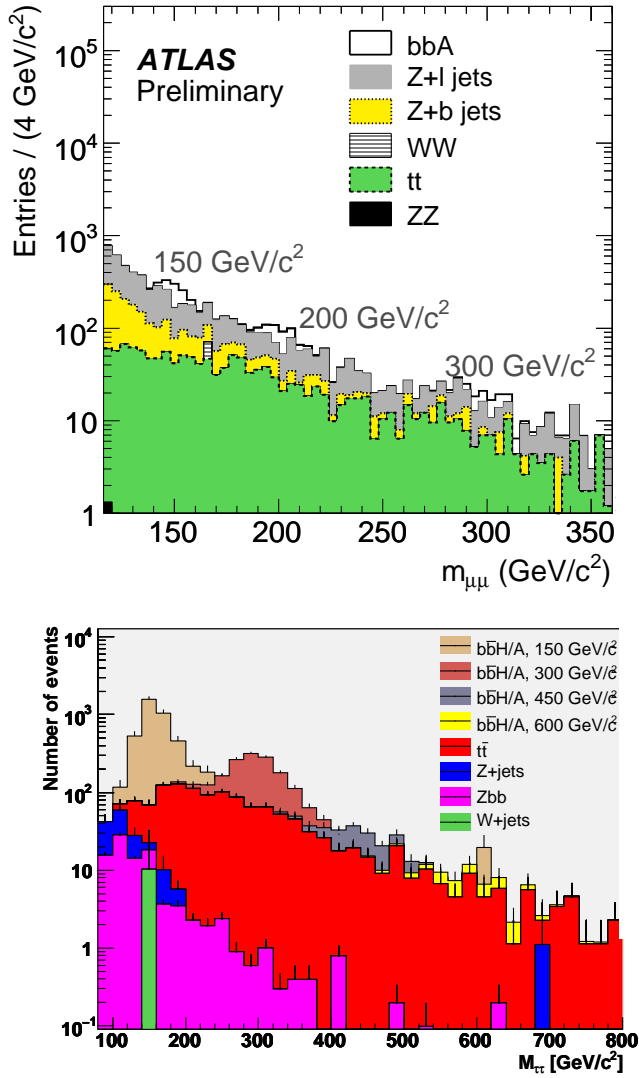


Figure 1.26: Invariant mass distributions of muon pairs in MSSM $A/H \rightarrow \mu^+\mu^-$ decays for Higgs masses m_A of 150, 200 and $300 \text{ GeV}/c^2$ (top picture) and of τ pairs in MSSM $A/H \rightarrow \tau^+\tau^-$ events for $m_A = 150, 300, 450$ and $600 \text{ GeV}/c^2$ together with the expected background contributions for $\tan\beta = 30$ and an integrated luminosity of 30 fb^{-1} .

constructed because of the undetected neutrinos in the final state, these events can only be distinguished as an excess of events with reconstructed τ leptons and large missing transverse energy above the high background of Standard model decays of top quark pairs and of $Z \rightarrow \tau\tau$ decays. In Fig. 1.28, the 5σ discovery region in the $\tan\beta$ - m_{H^\pm} plane for a charged Higgs boson in the above production and decay mode is shown for an integrated luminosity of 1 fb^{-1} assuming a systematic error on the background estimation of 10% [12]. We are developing methods for a reliable determination of

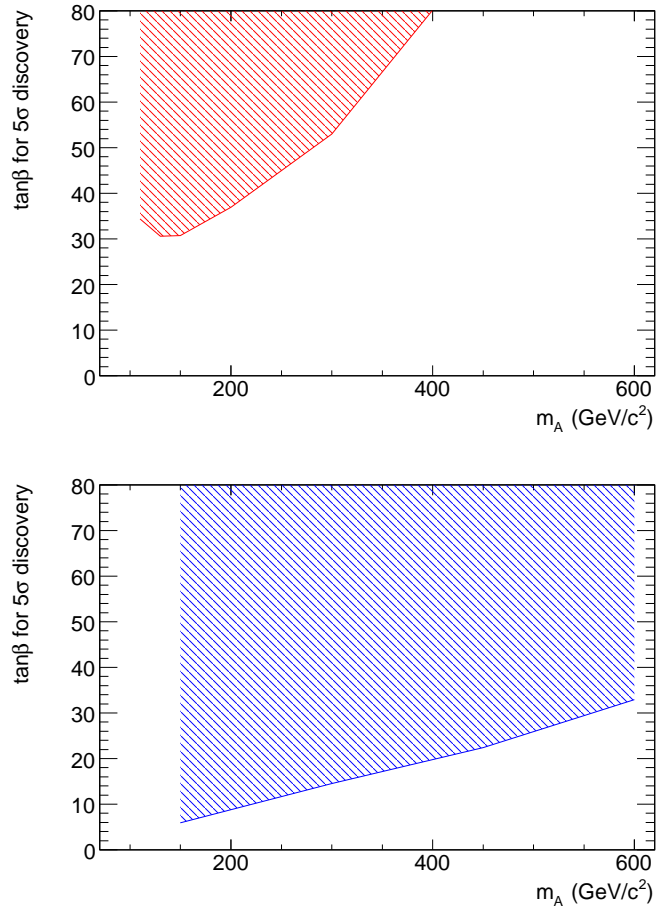


Figure 1.27: The regions of the $\tan\beta$ - m_A parameter space (hatched) in which $A/H/h \rightarrow \mu^+\mu^-$ (top picture) and $A/H \rightarrow \tau^+\tau^-$ decays (bottom picture) can be observed with at least 5σ significance for an integrated luminosity of 10 fb^{-1} .

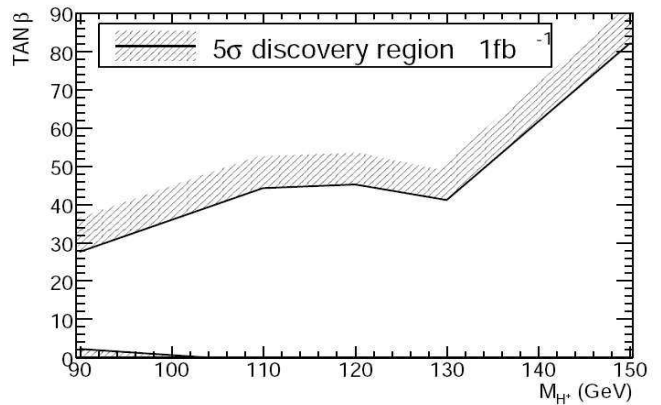


Figure 1.28: The region of the $\tan\beta$ - m_{H^\pm} parameter space (hatched region and above) in which charged Higgs decays $H^\pm \rightarrow \tau^\pm \nu_\tau$ can be observed with at least 5σ significance for an integrated luminosity of 1 fb^{-1} .

the background contributions from the LHC data using methods similar to the ones described in the next section.

1.2.5 Searches for Supersymmetric Particles

Supersymmetry (SUSY) is the most favoured extension of the Standard Model. The new symmetry uniting fermions and bosons predicts for each Standard Model particle a new supersymmetric partner with spin quantum number differing by $1/2$. Supersymmetry provides a natural explanation for Higgs boson masses near the electroweak scale if the supersymmetry breaking scale is not much higher. If, as expected in this case, the superpartners of the Standard Model particles, in particular squarks and gluinos (the superpartners of quarks and gluons with spin 0 and $1/2$, respectively) have masses below $1 \text{ TeV}/c^2$, they will be copiously produced at the LHC. Since details of the SUSY breaking mechanism are not known, possible supersymmetric extensions of the Standard Model populate a large parameter space preventing a prediction of the superpartner masses.

In supersymmetric models without R-parity violation, the superpartners eventually decay into the lightest supersymmetric particle which is stable and weakly interacting and escapes detection. Typical signatures for decays of supersymmetric particles in the ATLAS experiment are, therefore, large missing energy, energetic hadron jets and up to three leptons produced in the decay chain.

Our studies concentrate on the early discovery of such inclusive signatures of supersymmetric particle production and decays. The ATLAS detector is well equipped for such measurements. Reliable estimation of the background contributions from $t\bar{t}$, W/Z + jets and generic QCD jet production is the main concern. $t\bar{t}$ production is the dominant background for large missing energies where the signal is expected. With Monte Carlo simulations we have studied methods for determining the $t\bar{t}$ background from data [13]. First a pure background sample is selected using kinematic constraints specific for $t\bar{t}$ production. The properties of this control sample are then extrapolated to the SUSY signal region by replacing parts of the event topology by signatures specific for the signal events.

Fig. 1.29 shows the missing transverse energy distributions for supersymmetric particle decays with one (top picture) and with two leptons in the final state, respectively, for a representative SUSY model together with the expected $t\bar{t}$ background after accumulation of only

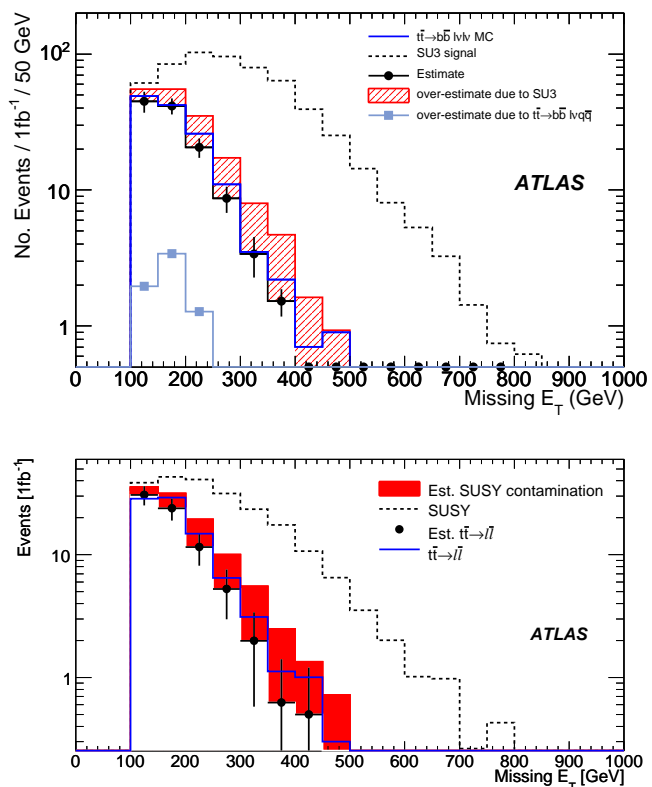


Figure 1.29: Missing transverse energy E_T^{miss} distributions for simulated supersymmetric particle decays with one (top picture) and with two leptons (bottom picture) in the final state in a representative SUSY model (SU3) together with the $t\bar{t}$ backgrounds estimated from a control data sample (see text) for an integrated luminosity of 1 fb^{-1} [13]. The uncertainties in the background estimates due to contamination of the control sample with signal events are indicated.

1 fb^{-1} of integrated luminosity. The uncertainties in the background estimates due to contamination of the control sample with signal events are also shown. For the model studied, a clear excess of SUSY events above the background is visible at large missing transverse energy.

In addition, we pursue searches for long-lived particles detectable only in the muon spectrometer and for lepton number violating τ decays into three muons which occur in various extensions of the Standard Model including many supersymmetric models.

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