# Testing the Standard Model of Elementary Particle Physics II

2nd lecture

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# **4. Recent experimental Tests on the Standard Model of Particle Physics**



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#### 4.1 Precision Measurements of the Electroweak Interaction



# The W boson

- Predicted in the 1968 by Sheldon Glashow, Abdus Salam und Steven Weinberg
- Discovered in 1983 by the UA1 and UA2 collaborations at the SppS (CERN)
- Properties:
  - $\circ$  m<sub>w</sub> = 80.385 ± 0.015 GeV
  - $\circ$   $\Gamma_{W} = 2.085 \pm 0.042 \text{ GeV}$
  - ο τ<sub>w</sub> ≈ 3 · 10<sup>-25</sup> s

| Decay Mode              | BR                   |  |  |  |  |  |
|-------------------------|----------------------|--|--|--|--|--|
| $W \rightarrow ev$      | $(10.71 \pm 0.16)\%$ |  |  |  |  |  |
| $W  ightarrow \mu \nu$  | $(10.63 \pm 0.15)\%$ |  |  |  |  |  |
| $W \to \tau \nu$        | $(11.38 \pm 0.21)\%$ |  |  |  |  |  |
| $W \rightarrow hadrons$ | $(67.41 \pm 0.27)\%$ |  |  |  |  |  |

- W boson decay (Lepton universality):
  - All three types of charged lepton particles interact in the same way with other particles.
  - The three lepton types are created equally often in particle transformations, or decays (once differences in their mass are accounted for)

# The Z boson

- Predicted in the 1960s by Sheldon Glashow, Abdus Salam und Steven Weinberg
- Discovered in 1983 by the UA1 and UA2 collaborations at the SppS (CERN)
- Properties:
  - $\circ$  m<sub>z</sub> = 91.1875 ± 0.0021 GeV
  - $\circ$   $\Gamma_{z}$  = 2.4952 ± 0.0023 GeV
  - $\circ$  T<sub>Z</sub> ≈ 2.5 · 10<sup>-25</sup> s

| Decay Mode                 | BR                      |  |  |  |  |
|----------------------------|-------------------------|--|--|--|--|
| $Z \rightarrow e^+ e^-$    | $(3.3632 \pm 0.0042)\%$ |  |  |  |  |
| $Z  ightarrow \mu^+ \mu^-$ | $(3.3662 \pm 0.0066)\%$ |  |  |  |  |
| $Z  ightarrow 	au^+ 	au^-$ | $(3.3696 \pm 0.0083)\%$ |  |  |  |  |
| $Z \rightarrow invisible$  | $(20.000 \pm 0.055)\%$  |  |  |  |  |
| $Z \rightarrow hadrons$    | $(69.911 \pm 0.056)\%$  |  |  |  |  |

#### 4.1.1 Measurements on the Z<sup>o</sup> boson resonance





- Between 1989 and 1995 (LEP phase I), the LEP experiments measured the production of fermion-antifermion pairs in the e<sup>+</sup>e<sup>-</sup> annihilation at the Z<sup>0</sup> peak (i.e. at a center-of-mass energy of √s = m<sub>7</sub>)
- At this energy, electrons and positrons create almost exclusively real (i.e. on-shell) Z<sup>0</sup> bosons
  - Basically at rest
  - With subsequent decay into fermion-antifermion pairs  $(2m_f \le m_7)$
- In general, significant contributions of interference effects between photon- and Z<sup>0</sup> exchange (electroweak interference between these two NC-processes):

• The fermion-antifermion pair production cross section (via  $e^+e^-$  annihilation) is a function of the center-of-mass energy  $\sqrt{s} = 2E_e$  and the scattering angle  $\theta$  and is given (in lowest order) via:

$$\begin{aligned} \frac{d\sigma}{d\Omega} \left( e^+ e^- \to f\bar{f} \right) &= \operatorname{const.} \cdot s \cdot \left| \mathcal{M}_{QED} + \mathcal{M}_{weak} \right|^2 \\ &= \operatorname{const.} \cdot s \cdot \left[ \left| \mathcal{M}_{QED} \right|^2 + 2\mathcal{R}e \left( \mathcal{M}_{QED} \mathcal{M}_{weak} \right) + \left| \mathcal{M}_{weak} \right|^2 \right] \\ &= N_C^f \frac{\alpha^2}{4s} \left[ C_1^f (1 + \cos^2 \theta) + C_2^f \cos \theta \right] \\ &= \frac{1}{2\pi} \sigma_{f\bar{f}} \left[ \frac{3}{8} \left( 1 + \cos^2 \theta \right) + A_{FB}^f \cos \theta \right] \end{aligned}$$

with:

$$C_1^f(s) = Q_e^2 Q_f^2 + 8Q_e Q_f v_e v_f \mathcal{R}e(\chi(s)) + 16(v_e^2 + a_e^2)(v_f^2 + a_f^2)|\chi(s)|^2$$

$$C_2^f(s) = 16Q_eQ_fa_ea_f\mathcal{R}e(\chi(s)) + 32v_ea_ev_fa_f|\chi(s)|^2$$

$$\sigma_{f\bar{f}}(s) = \int_0^{2\pi} \int_{-1}^{+1} \frac{d\sigma}{d\Omega} d\cos\theta d\Phi = N_C^f \frac{4\pi\alpha^2}{3s} \cdot C_1^f(s)$$

$$\begin{aligned} A_{FB}^{f}(s) &= \frac{\int_{0}^{+1} \frac{d\sigma}{d\Omega} d\cos\theta - \int_{-1}^{0} \frac{d\sigma}{d\Omega} d\cos\theta}{\int_{-1}^{+1} \frac{d\sigma}{d\Omega} d\cos\theta} \\ &= \frac{N_{F} - N_{B}}{N_{F} + N_{B}} = \frac{3}{8} \cdot \frac{C_{2}^{f}(s)}{C_{1}^{f}(s)} \end{aligned}$$

Coupling constants of the weak vector- and axial-currents v<sub>i</sub> & a<sub>i</sub>

 $Z^0$  boson propagator  $\chi(s)$ 

Forward-Backward asymmetry A<sub>FB</sub>

Coupling constants of the weak vector- and axial-currents:

$$v_f = I_f^0 - 2Q_f \sin^2 \theta_W \equiv \frac{1}{2} - 2Q_f \sin^2 \theta_W$$
  
$$a_f = I_f^0 \equiv \frac{1}{2}$$
  
In the center-of-mass system:  
$$q^2 = (p_1 + p_2)^2 =: s \text{ and } \sqrt{s} = 2E_e$$

Z<sup>0</sup> boson propagator:

 $G_F = 1.166 3787 \times 10^{-5} \text{ GeV}^{-2}$  (Fermi coupling constant)

$$\chi(s) = \frac{1}{16\sin^2\theta_W \cos^2\theta_W} \cdot \frac{s}{s - M_z^2 + iM_z\Gamma_z}$$
$$= \frac{G_F M_z^2}{8\pi\alpha\sqrt{2}} \cdot \frac{s}{s - M_z^2 + iM_z\Gamma_z}$$

• The QED contribution (only photon exchange) is described via:

$$\frac{d\sigma}{d\Omega}(e^+e^- \to f\bar{f}) = N_C^f \frac{\alpha^2 Q_f^2}{4s} (1 + \cos^2\theta)$$
$$\sigma_{f\bar{f}} = \int \frac{d\sigma}{d\Omega} d\Omega = \frac{4\pi\alpha^2}{3s} = \frac{87 \text{ nb}}{s[\text{GeV}^2]} \cdot Q_f^2 \cdot N_C^f$$

For leptons  $N_c$  is 1, while for quarks it is equal to 3.

• Thus:

$$R = \frac{\sigma\left(e^+e^- \to \sum_{q(E_{CMS} > 2m_q)} q\bar{q}\right)}{\sigma(e^+e^- \to \mu^+\mu^-)}$$
$$= 3 \cdot \sum_{q(E_{CMS} > 2m_q)} Q_q^2$$

- Since the Z<sup>0</sup> boson has a finite lifetime τ<sub>z</sub> and width Γ<sub>z</sub>, the curve of the Z<sup>0</sup> resonance is described by the Breit-Wigner function
- At the **pole mass** of the Z<sup>0</sup> boson, the cross section is given via:

$$\begin{aligned} \sigma_{f\bar{f}}^{0}(s=M_{Z}^{2}) &= N_{C}^{f}\frac{G_{F}^{2}M_{Z}^{4}}{6\pi\Gamma_{Z}^{2}}\cdot(v_{e}^{2}+a_{e}^{2})(v_{f}^{2}+a_{f}^{2}) \\ &= \frac{12\pi\Gamma_{Z}^{e}\Gamma_{Z}^{f}}{M_{Z}^{2}\Gamma_{Z}^{2}} \end{aligned}$$

• Decay width of the Z<sup>0</sup> boson are given via:

$$\begin{split} \Gamma(Z^0 \to f\bar{f}) &\equiv \Gamma_Z^f = \frac{G_F M_Z^3}{6\pi\sqrt{2}} (v_f^2 + a_f^2) \cdot N_C^f \\ \Gamma(Z^0 \to e^+ e^-) &\equiv \Gamma_Z^e = \frac{G_F M_Z^3}{6\pi\sqrt{2}} (v_e^2 + a_e^2) \\ \Gamma_Z &\equiv \sum_f \Gamma_Z^f \\ \Gamma_Z^{had} &\equiv \sum_q \Gamma_Z^f \\ A_{FB}^f = \frac{3}{4} A_e A_f \quad \text{with} \quad A_f = \frac{2v_f a_f}{v_f^2 + a_f^2} \end{split}$$











| Number of Events |                               |      |      |      |                               |     |     |     |     |      |
|------------------|-------------------------------|------|------|------|-------------------------------|-----|-----|-----|-----|------|
|                  | $Z \rightarrow q\overline{q}$ |      |      |      | $Z \rightarrow \ell^+ \ell^-$ |     |     |     |     |      |
| Year             | A                             | D    | L    | 0    | LEP                           | A   | D   | L   | 0   | LEP  |
| 1990/91          | 433                           | 357  | 416  | 454  | 1660                          | 53  | 36  | 39  | 58  | 186  |
| 1992             | 633                           | 697  | 678  | 733  | 2741                          | 77  | 70  | 59  | 88  | 294  |
| 1993             | 630                           | 682  | 646  | 649  | 2607                          | 78  | 75  | 64  | 79  | 296  |
| 1994             | 1640                          | 1310 | 1359 | 1601 | 5910                          | 202 | 137 | 127 | 191 | 657  |
| 1995             | 735                           | 659  | 526  | 659  | 2579                          | 90  | 66  | 54  | 81  | 291  |
| Total            | 4071                          | 3705 | 3625 | 4096 | 15497                         | 500 | 384 | 343 | 497 | 1724 |

The qq and  $l^+l^-$  event statistics, in units of  $10^3$ , used for Z analyses by the experiments ALEPH (A), DELPHI (D), L3 (L) and OPAL (O).

- Angular distribution observed in the lepton pair production at LEP
  - Characteristic forward background asymmetry A<sub>FB</sub>

 $\sqrt{s}$  : centre-of-mass energy

 $\forall \text{s'}: \text{effective centre-of-mass energy after}$  initial-state photon radiation



#### Hadronic cross section



# Hadronic cross section

 Measurement of number of light neutrinos (2m<sub>v</sub> < m<sub>z</sub>):

$$N_{
u} = 2.9841 \pm 0.0083$$

- Equivalent to a measurement of quark-lepton generations
- Measured in 1989 by the LEP experiments (but also by MARK-II at SLC)



# Hadronic cross section

- The total decay width of the Z<sup>0</sup> boson Γ<sub>z</sub> depends on the number of neutrino generations
  - $\circ$  while  $\Gamma^{had}$  does not
- $\rightarrow$  The hadronic cross section

$$\sigma_{\rm had}^{0} = \frac{12\pi\Gamma_Z^e\Gamma_Z^{\rm had}}{M_Z^2\Gamma_Z^2}$$

varies as a function of the neutrino numbers



# 4.1.2 Measurement of the W boson production process



### W boson production at LEP

- Between 1996 and 2000, the center-of-mass energy of the LEP collider was increased (step-by-step) to:
  - 130, 161, 172 GeV (1996)
  - 183 GeV (1997)
  - 189 GeV (1998)
  - 192, 196, 200, 202 GeV (1999)
  - 208 GeV (2000)
- With these scans, the W<sup>+</sup>W<sup>-</sup> production threshold (and thus the W boson mass) was measured precisely

#### W boson production (and decay) at LEP





### W boson mass measurements (at ALEPH)

- The W boson mass and width are extracted by fitting simulated invariant mass spectra to the observed distributions.
  - Use an unbinned maximum likelihood procedure





Taken from: https://arxiv.org/pdf/hep-ex/0605011.pdf

#### W boson mass measurements (at LEP & Tevatron)











Width of the W Boson



- Electroweak radiative corrections modify the calculation of Z-pole observables.
  - Virtual W and Higgs bosons as well as top quarks contribute via loops
- The contributions from radiative corrections depend on the masses of these particles:
  - $\circ$  m<sub>W</sub> and m<sub>top</sub> (quadratic)
  - m<sub>H</sub> (logarithmic)

Taken from: https://arxiv.org/pdf/hep-ex/0509008.pdf





Higher-order corrections to the gauge boson propagators due to boson loops

- Comparisons with the high precision measurements allow for an indirect determination of  $m_{W}^{}$ ,  $m_{top}^{}$  and  $m_{H}^{}$ 
  - **Prediction:** 
    - m<sub>w</sub> = 80.364 ± 0.021 GeV
    - m<sub>top</sub> = 172.3 ± 10 GeV
    - m<sub>H</sub> = 91 ± 45 GeV
  - Direct measurement:
    - m<sub>w</sub> = 80.399 ± 0.025 GeV
    - m<sub>top</sub> = 173.2 ± 0.8 GeV (early: TEVATRON + LHC combination)
    - m<sub>H</sub> = 125.09 ± 0.21(stat) ± 0.11 (syst) GeV



The shaded area denotes the indirect determination of the top quark mass at the 68% Confidence Level (C.L.) derived from the analysis of radiative corrections within the framework of the SM using precision electroweak measurements.



Taken from <a href="https://arxiv.org/pdf/hep-ex/0509008.pdf">https://arxiv.org/pdf/hep-ex/0509008.pdf</a>39

#### 4.1.4 W & Z bosons at the LHC



#### W/Z + jets production



#### **Charge asymmetry in W boson production**

• Parton distribution functions (PDFs) of u and d quarks in the proton differ (mainly due to the fact that protons contain two valence u quarks and one valence d quark)





## W boson production

- W boson candidate events are selected by requiring:
  - Exactly one identified electron or muon
  - MET > 25 GeV
  - $\circ$  m<sub>T</sub> > 50 GeV
- Roughly 20% of all selected events stemm from background processes:
  - Most dominant contributions:
    - Multijets (10%)
    - $\blacksquare \quad Z \to \ell \ell \ (5\%)$





# Z boson production

- Z boson candidate events are selected by requiring:
  - Exactly two leptons:
    - Same flavour & opposite charge
    - 66 < m<sub>ℓℓ</sub> < 116 GeV
- Roughly 0.5% of all selected events stemm from background processes:





Run: 267639 Event: 173263110 2015-06-14 13:13:03 CEST

 $Z \rightarrow e^+e^-$ 

proton-proton collisions at 13 TeV





### **W/Z production cross sections**

From: https://arxiv.org/abs/1603.09222



W & Z boson production cross section measurements are sensitive to the PDF sets  $\rightarrow$  Can constrain them

#### **EW Z boson production cross section**

- Insights into the mechanism of electroweak (EW) symmetry breaking can be achieved via (VBS) processes
  - Via studies of vector boson self interaction
- Sensitive to BSM contributions:
  - Heavy-vector triplets
  - Extra dimensions





$$\sigma_{\rm EW}$$
 = 37.4 ± 3.5 (stat) ± 5.5 (syst) fb

Taken from: https://arxiv.org/pdf/2006.15458.pdf

#### **Cross section for a** *Z* **boson produced in association with** *b***-jets**



50



#### **Diboson production (WW, WZ, ZZ)**







Run Number: 284420, Event Number: 546213887

Date: 2015-11-02 00:56:41 CET



#### WZ production cross section measurements





# **Same-sign WW boson pair production**

- Insights into the mechanism of electroweak (EW) symmetry breaking can be achieved via vector boson scattering (VBS) processes
  - Via studies of vector boson self interaction
- Sensitive to anomalous quartic gauge couplings (aQGC)
- Same-sign W<sup>±</sup>W<sup>±</sup> channel is promising due to small background yields from SM processes





 $\sigma$  BR(W<sup>±</sup>W<sup>±</sup> $\rightarrow$   $\ell^{\pm}\ell^{\pm}$ ) = 3.98 ± 0.37 (stat) ± 0.25 (syst) fb

From https://arxiv.org/pdf/2005.01173.pdf

#### **Triboson production**

 $V_1$ 

 $V_2$ 

- Rare processes with cross sections in the order of 1pb
- Important background to Di-Higgs searches
- Test of the non-Abelian gauge structure of the SM
  - Any deviations from the SM prediction would provide hints of new physics at higher energy scales

